

# Particular Vibration Spectrum of Antimony Trichloride Caused by Bond Fluctuation

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Vibration spectra of solid and molten antimony trichloride were recorded. In the solid state, vibration frequencies in addition to those of the  $\text{SbCl}_3$  molecules with their  $C_{3v}$  symmetry indicate bond fluctuations according to  $2 \text{SbCl}_3 \rightleftharpoons \text{Sb}^+ + \text{SbCl}_6^-$ .

**Key words:**  $\text{SbCl}_3$ , Vibration spectrum, Bond fluctuation, Electron transfer.

## Introduction

In the past, antimony trichloride has extensively been investigated by structure sensitive methods, especially by vibration spectroscopy [1–8].  $\text{SbCl}_3$  crystallizes in the space group  $\text{Pnma}$  ( $D_{2h}^{16}$ ), and it is isotopic with several trihalides of Group V elements like  $\text{NCl}_3$ ,  $\text{PCl}_3$ ,  $\text{PBr}_3$ ,  $\text{AsCl}_3$ ,  $\beta\text{-SbBr}_3$  and  $\text{BiCl}_3$ .

In the solid state, the vibration spectra of these compounds differ considerably from their melt and gas spectra. The last ones can be explained by the  $C_{3v}$  symmetry of the  $\text{AB}_3$  molecule.

The more complex vibration spectra of the solids are mostly caused by intermolecular interactions. Especially the Raman spectra of the phosphorus and arsenic trihalides exhibit many additional bands [9–10] while  $\text{SbCl}_3$ ,  $\text{SbBr}_3$ ,  $\text{BiCl}_3$  and  $\text{BiBr}_3$  show a simpler splitting pattern [3–8, 11]. For  $\text{PBr}_3$  below 200 K ten Raman active deformation and stretching frequencies were observed while for  $\text{SbCl}_3$  below 340 K there are only seven Raman bands.

Obviously different interactions exist in these solids despite the same crystal structure. This has not been sufficiently considered in the cited investigations, so that a new discussion of the  $\text{SbCl}_3$  vibration spectrum is justified.

## Results

The FIR spectra as well as the Raman spectrum of crystalline  $\text{SbCl}_3$  along with the Raman spectrum of a solidified melt sample of  $\text{SbCl}_3$  with 15.0 mol%  $\text{RbCl}$  are presented in Figure 1. Table 1 summarizes the observed frequencies ( $\text{cm}^{-1}$ ) of solid  $\text{SbCl}_3$  and the  $\text{SbCl}_3\text{-RbCl}$  mixture along with their intensities and literature data [3, 15]. Especially the low temperature FIR spectrum of  $\text{SbCl}_3$  at 8 K demonstrates that it is much more complex than formerly assumed.

## Description of the Structure

Based on data of Lindquist and Niggli [1] a section of the crystal structure of  $\text{SbCl}_3$  is sketched in Figure 2. The unit cell with space group  $\text{Pnma}$  consists of four molecular units. In Table 2 the numbers and distances of the  $\text{Sb-Cl}$  coordinations are given. Therefore Lipka [2] proposed a 3-dimensional structure built up by bicapped-trigonal prisms, so that a coordination number 8 for antimony results (Figure 2).

In general, it is difficult to make predictions for real chemical bonds only on the basis of X-ray structure data. So the  $\text{Sb-Cl}$  coordinations with distances in the range of 345–400 pm cannot be considered as distinct chemical bonds.

## Experimental

Spectroscopically pure  $\text{SbCl}_3$  was prepared by the reaction of antimony pentachloride and antimony according to



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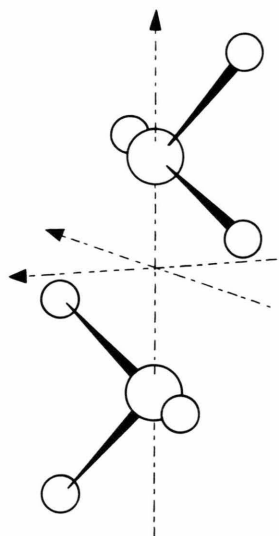


Fig. 3. Vibration spectroscopically relevant unit cell with two SbCl<sub>3</sub> units and the factor group C<sub>2h</sub>.

sample compartment have already been described [13, 14]. The room temperature far infrared (FIR) spectrum was recorded with a modified Beckman interferometer FS 720 interfaced to a Kontron KAP 1000 computer for acquiring and manipulating the data, and the low temperature FIR spectrum with a Bruker IFS 113 FT spectrometer.

## Discussion

So far the interpretation of the SbCl<sub>3</sub> vibration spectrum is based on factor group analyses with four SbCl<sub>3</sub> molecules integrated in the unit cell [4]. The four fundamentals of the SbCl<sub>3</sub> molecules (2A<sub>1</sub> + 2E) then split into 24 vibrations according to

$$\begin{aligned} \Gamma_{\text{vib}}(D_{2h}) = & 4 A_g(\text{RE}) + 2 B_{1g}(\text{RE}) + 4 B_{2g}(\text{RE}) \\ & + 2 B_{3g}(\text{RE}) + 2 A_u + 4 B_{1u}(\text{IR}) \\ & + 2 B_{2u}(\text{IR}) + 4 B_{3u}(\text{IR}). \end{aligned} \quad (2)$$

Species in A<sub>u</sub> are inactive and the mutual exclusion rule is valid. Instead of the twelve Raman active deformations and stretchings expected only seven could be observed.

With an equalization of the crystallographic with the spectroscopic unit cell, 3-dimensional interactions are assumed which at least are so high that they deliver a contribution to the vibrational behavior. If one of these contributions exceeds the others significantly,

Table 3. Splitting schematic for crystalline SbCl<sub>3</sub> and proposed assignment of the vibration frequencies (cm<sup>-1</sup>).

Point group	Site group	Factor group	Factor group	$\nu/\text{cm}^{-1}$
C <sub>3v</sub>	C <sub>s</sub>	D <sub>2h</sub> <sup>1)</sup>	C <sub>2h</sub> <sup>2)</sup>	
$\nu_1 (A_1)$	A'	A <sub>g</sub> B <sub>2g</sub> B <sub>1u</sub> B <sub>3u</sub>	A <sub>g</sub> B <sub>u</sub>	342 354
$\nu_2 (E)$	A'	A <sub>g</sub> B <sub>2g</sub> B <sub>1u</sub> B <sub>3u</sub>	A <sub>g</sub> B <sub>u</sub>	318 363
	A''	B <sub>1g</sub> B <sub>3g</sub> A <sub>u</sub> B <sub>2u</sub>	B <sub>g</sub> A <sub>u</sub>	313 348
$\nu_3 (A_1)$	A'	A <sub>g</sub> B <sub>2g</sub> B <sub>1u</sub> B <sub>3u</sub>	A <sub>g</sub> B <sub>u</sub>	151 159
$\nu_4 (E)$	A'	A <sub>g</sub> B <sub>2g</sub> B <sub>1u</sub> B <sub>3u</sub>	A <sub>g</sub> B <sub>u</sub>	139 135
	A''	B <sub>1g</sub> B <sub>3g</sub> A <sub>u</sub> B <sub>2u</sub>	B <sub>g</sub> A <sub>u</sub>	135 141

then vibrations move close together and cannot be separated any more, and a distinct assignment of the vibration modes is impossible.

In this case it is suitable and allowed to reduce the basic unit of a factor group analysis, that is to halve the crystallographic unit cell. As a result, the factor group C<sub>2h</sub> (Figure 3) is obtained. Thus, the SbCl<sub>3</sub> structure may be imagined as a linear chain of AB<sub>3</sub> molecules, and according to Table 2 only the intermolecular Sb–Cl distances to 346 pm are considered as interactions; therefore the coordination number of antimony is five.

A vibrational analysis with two SbCl<sub>3</sub> molecules in the unit cell leads to the following result: (3)

$$\Gamma_{\text{vib}}(C_{2h}) = 4 A_g(\text{RE}) + 2 B_g(\text{RE}) + 2 A_u(\text{IR}) + 4 B_u(\text{IR}).$$

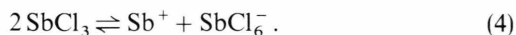
As expected, the number of the vibrations is halved, and the Raman spectrum should show three deformation as well as three stretching frequencies.

In Table 3 the splitting schematic for the different models used and the assignment of the frequencies for the reduced spectroscopic unit cell are summarized and compared. These results demonstrate that the number of species in the deformation and stretching region must always have the same value independently of the ansatz chosen. This is contrary to the experimental Raman spectrum (Fig. 1d) where four deformation and three stretching vibrations are observed. Therefore the Raman band at 170 cm<sup>-1</sup> in the SbCl<sub>3</sub> Raman spectrum is *not* caused by the SbCl<sub>3</sub> molecules and/or its intermolecular interactions in the lattice.

Obviously, small additions of alkali chloride disturb considerably the SbCl<sub>3</sub> structure (cf. Fig. 1c; SbCl<sub>3</sub> + 15 mol% RbCl). Here the Raman bands at 318 and 139 cm<sup>-1</sup> are no more split. Nevertheless, the Raman vibration at 170 cm<sup>-1</sup> remains with unchanged intensity.

In the SbCl<sub>3</sub> crystal structure rapid electron transfer or bond fluctuations are imaginable, which will cause structural defects, and could be a possible reason for the existence of the Raman band at 170 cm<sup>-1</sup>. It is striking that the complex anion SbCl<sub>6</sub><sup>-</sup> shows a strong Raman deformation at 170 cm<sup>-1</sup> (ν<sub>5</sub>(F<sub>2g</sub>), too, which has been observed in several compounds of the type Me<sup>I</sup>Sb<sup>V</sup>Cl<sub>6</sub> (Me<sup>I</sup> = K<sup>+</sup>, Rb<sup>+</sup>, Cs<sup>+</sup>, NO<sup>+</sup>) [16–20] as well as in the mixed-valent type Me<sub>4</sub><sup>I</sup>Sb<sup>III</sup>Cl<sub>6</sub>Sb<sup>V</sup>Cl<sub>6</sub> [21–24]. Possibly SbCl<sub>6</sub><sup>-</sup> may originate by a rapid

electron or bond fluctuation in the SbCl<sub>3</sub> crystal structure according to a more formal equilibrium reaction



The ν<sub>1</sub>(A<sub>g</sub>) Raman stretching frequency of the SbCl<sub>6</sub><sup>-</sup> anion at 340 cm<sup>-1</sup> could be covered by an SbCl<sub>3</sub> fundamental. Normally the ν<sub>2</sub>(E<sub>g</sub>) Raman band of the octahedral SbCl<sub>6</sub><sup>-</sup> at 284 cm<sup>-1</sup> is very weak, so that the lack of such a band in the SbCl<sub>3</sub> spectrum could be explained.

As a further hint for the existence of SbCl<sub>6</sub><sup>-</sup> in solid SbCl<sub>3</sub> the low temperature FIR spectrum (Fig. 1a) may be considered: The absorption bands at 333 and 189 cm<sup>-1</sup> correspond to SbCl<sub>6</sub><sup>-</sup> stretching and deformation frequencies (ν<sub>3</sub>, ν<sub>4</sub>(F<sub>1u</sub>)).

The great instability of Sb(I) compounds [25], whose existence in the condensed state was not proved distinctly until now, should cause a short life time. Possibly this could explain that an expected isomery shift for solid SbCl<sub>3</sub> is not found in Mößbauer spectra (Sb(I), Sb(III), Sb(V)) [26].

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